

# RESULTS FROM ANALYTIC FUNCTION THEORY

## C.1 Introduction

This appendix summarizes key results from analytic function theory leading to the Cauchy Integral formula and its consequence, the Poisson–Jensen formula.

## C.2 Independence of Path

Consider functions of two independent variables,  $x$  and  $y$ . (The reader can think of  $x$  as the real axis and  $y$  as the imaginary axis.)

Let  $P(x, y)$  and  $Q(x, y)$  be two functions of  $x$  and  $y$ , continuous in some domain  $D$ . Say we have a curve  $C$  in  $D$ , described by the parametric equations

$$x = f_1(t), \quad y = f_2(t) \quad (\text{C.2.1})$$

We can then define the following line integrals along the path  $C$  from point  $A$  to point  $B$  inside  $D$ .

$$\int_A^B P(x, y) dx = \int_{t_1}^{t_2} P(f_1(t), f_2(t)) \frac{df_1(t)}{dt} dt \quad (\text{C.2.2})$$

$$\int_A^B Q(x, y) dy = \int_{t_1}^{t_2} Q(f_1(t), f_2(t)) \frac{df_2(t)}{dt} dt \quad (\text{C.2.3})$$

**Definition C.1.** *The line integral  $\int P dx + Q dy$  is said to be **independent of the path** in  $D$  if, for every pair of points  $A$  and  $B$  in  $D$ , the value of the integral is independent of the path followed from  $A$  to  $B$ .*

□□□

We then have the following result.

**Theorem C.1.** *If  $\int Pdx + Qdy$  is independent of the path in  $D$ , then there exists a function  $F(x, y)$  in  $D$  such that*

$$\frac{\partial F}{\partial x} = P(x, y); \quad \frac{\partial F}{\partial y} = Q(x, y) \quad (\text{C.2.4})$$

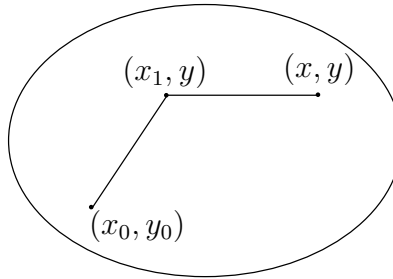
*hold throughout  $D$ . Conversely, if a function  $F(x, y)$  can be found such that (C.2.4) hold, then  $\int Pdx + Qdy$  is independent of the path.*

**Proof**

Suppose that the integral is independent of the path in  $D$ . Then, choose a point  $(x_0, y_0)$  in  $D$  and let  $F(x, y)$  be defined as follows

$$F(x, y) = \int_{x_0, y_0}^{x, y} Pdx + Qdy \quad (\text{C.2.5})$$

where the integral is taken on an arbitrary path in  $D$  joining  $(x_0, y_0)$  and  $(x, y)$ . Because the integral is independent of the path, the integral does indeed depend only on  $(x, y)$  and defines the function  $F(x, y)$ . It remains to establish (C.2.4).



**Figure C.1.** Integration path

For a particular  $(x, y)$  in  $D$ , choose  $(x_1, y)$  so that  $x_1 \neq x$  and so that the line segment from  $(x_1, y)$  to  $(x, y)$  in  $D$  is as shown in Figure C.1. Because of independence of the path,

$$F(x, y) = \int_{x_0, y_0}^{x_1, y} (Pdx + Qdy) + \int_{x_1, y}^{x, y} (Pdx + Qdy) \quad (\text{C.2.6})$$

We think of  $x_1$  and  $y$  as being fixed while  $(x, y)$  may vary along the horizontal line segment. Thus  $F(x, y)$  is being considered as function of  $x$ . The first integral on the right-hand side of (C.2.6) is then independent of  $x$ .

Hence, for fixed  $y$ , we can write

$$F(x, y) = \text{constant} + \int_{x_1}^x P(x, y) dx \quad (\text{C.2.7})$$

The fundamental theorem of Calculus now gives

$$\frac{\partial F}{\partial x} = P(x, y) \quad (\text{C.2.8})$$

A similar argument shows that

$$\frac{\partial F}{\partial y} = Q(x, y) \quad (\text{C.2.9})$$

Conversely, let (C.2.4) hold for some  $F$ . Then, with  $t$  as a parameter,

$$F(x, y) = \int_{x_1, y_1}^{x_2, y_2} P dx + Q dy = \int_{t_1}^{t_2} \left( \frac{\partial F}{\partial x} \frac{dx}{dt} + \frac{\partial F}{\partial y} \frac{dy}{dt} \right) dt \quad (\text{C.2.10})$$

$$= \int_{t_1}^{t_2} \frac{\partial F}{\partial t} dt \quad (\text{C.2.11})$$

$$= F(x_2, y_2) - F(x_1, y_1) \quad (\text{C.2.12})$$

□□□

**Theorem C.2.** *If the integral  $\int P dx + Q dy$  is independent of the path in  $D$ , then*

$$\oint P dx + Q dy = 0 \quad (\text{C.2.13})$$

*on every closed path in  $D$ . Conversely if (C.2.13) holds for every simple closed path in  $D$ , then  $\int P dx + Q dy$  is independent of the path in  $D$ .*

**Proof**

Suppose that the integral is independent of the path. Let  $C$  be a simple closed path in  $D$ , and divide  $C$  into arcs  $\vec{AB}$  and  $\vec{BA}$  as in Figure C.2.

$$\oint_C (P dx + Q dy) = \int_{AB} P dx + Q dy + \int_{BA} P dx + Q dy \quad (\text{C.2.14})$$

$$= \int_{AB} P dx + Q dy - \int_{AB} P dx + Q dy \quad (\text{C.2.15})$$

The converse result is established by reversing the above argument.

□□□

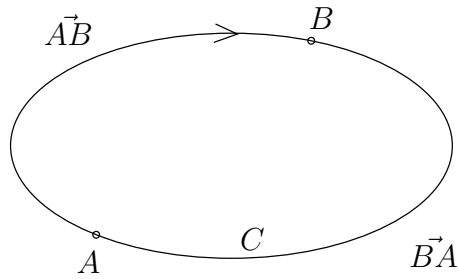


Figure C.2. Integration path

**Theorem C.3.** *If  $P(x, y)$  and  $Q(x, y)$  have continuous partial derivatives in  $D$  and  $\int Pdx + Qdy$  is independent of the path in  $D$ , then*

$$\frac{\partial P}{\partial y} = \frac{\partial Q}{\partial x} \quad \text{in } D \quad (\text{C.2.16})$$

**Proof**

By Theorem C.1, there exists a function  $F$  such that (C.2.4) holds. Equation (C.2.16) follows by partial differentiation.

□□□

Actually, we will be particularly interested in the converse to Theorem C.3. However, this holds under slightly more restrictive assumptions, namely a simply connected domain.

### C.3 Simply Connected Domains

Roughly speaking, a domain  $D$  is simply connected if it has no holes. More precisely,  $D$  is simply connected if, for every simple closed curve  $C$  in  $D$ , the region  $R$  enclosed by  $C$  lies wholly in  $D$ . For simply connected domains we have the following:

**Theorem C.4 (Green's theorem).** *Let  $D$  be a simply connected domain, and let  $C$  be a piecewise-smooth simple closed curve in  $D$ . Let  $P(x, y)$  and  $Q(x, y)$  be functions that are continuous and that have continuous first partial derivatives in  $D$ . Then*

$$\oint (Pdx + Qdy) = \int \int_R \left( \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dx dy \quad (\text{C.3.1})$$

where  $R$  is the region bounded by  $C$ .

**Proof**

We first consider a simple case in which  $R$  is representable in both of the forms:

$$f_1(x) \leq f_2(x) \quad \text{for } a \leq x \leq b \quad (\text{C.3.2})$$

$$g_1(y) \leq g_2(y) \quad \text{for } c \leq y \leq d \quad (\text{C.3.3})$$

Then

$$\int \int_R \frac{\partial P}{\partial y} dx dy = \int_a^b \int_{f_1(x)}^{f_2(x)} \frac{\partial P}{\partial y} dx dy \quad (\text{C.3.4})$$

One can now integrate to achieve

$$\int \int_R \frac{\partial P}{\partial y} dx dy = \int_a^b [P(x, f_2(x)) - P(x, f_1(x))] dx \quad (\text{C.3.5})$$

$$= \int_a^b P(x, f_2(x)) dx - \int_a^b P(x, f_1(x)) dx \quad (\text{C.3.6})$$

$$= \oint_C P(x, y) dx \quad (\text{C.3.7})$$

By a similar argument,

$$\int \int_R \frac{\partial Q}{\partial x} dx dy = \oint_C Q(x, y) dy \quad (\text{C.3.8})$$

For more complex regions, we decompose into simple regions as above. The result then follows.  $\square\square\square$

We then have the following converse to Theorem C.3.

**Theorem C.5.** *Let  $P(x, y)$  and  $Q(x, y)$  have continuous derivatives in  $D$  and let  $D$  be simply connected. If  $\frac{\partial P}{\partial y} = \frac{\partial Q}{\partial x}$ , then  $\oint P dx + Q dy$  is independent of path in  $D$ .*

**Proof**

Suppose that

$$\frac{\partial P}{\partial y} = \frac{\partial Q}{\partial x} \quad \text{in } D \quad (\text{C.3.9})$$

Then, by Green's Theorem (Theorem C.4),

$$\oint_c Pdx + Qdy = \int \int_R \left( \frac{\partial Q}{\partial x} - \frac{\partial P}{\partial y} \right) dxdy = 0 \quad (\text{C.3.10})$$

□□□

## C.4 Functions of a Complex Variable

In the sequel, we will let  $z = x + jy$  denote a complex variable. Note that  $z$  is not the argument in the Z-transform, as used at other points in the book. Also, a function  $f(z)$  of a complex variable is equivalent to a function  $\bar{f}(x, y)$ . This will have real and imaginary parts  $u(x, y)$  and  $v(x, y)$  respectively.

We can thus write

$$f(z) = u(x, y) + jv(x, y) \quad (\text{C.4.1})$$

Note that we also have

$$\begin{aligned} \int_C f(z)dz &= \int_C (u(x, y) + jv(x, y))(dx + jdy) \\ &= \int_C u(x, y)dx - \int_C v(x, y)dy + j \left\{ \int_C u(x, y)dy + \int_C v(x, y)dx \right\} \end{aligned}$$

We then see that the previous results are immediately applicable to the real and imaginary parts of integrals of this type.

## C.5 Derivatives and Differentials

Let  $w = f(z)$  be a given complex function of the complex variable  $z$ . Then  $w$  is said to have a derivative at  $z_0$  if

$$\lim_{\Delta z \rightarrow 0} \frac{f(z_0 + \Delta z) - f(z_0)}{\Delta z} \quad (\text{C.5.1})$$

exists and is independent of the direction of  $\Delta z$ . We denote this limit, when it exists, by  $f'(z_0)$ .

## C.6 Analytic Functions

**Definition C.2.** A function  $f(z)$  is said to be analytic in a domain  $D$  if  $f$  has a continuous derivative in  $D$ .

□□□

**Theorem C.6.** *If  $w = f(z) = u + jv$  is analytic in  $D$ , then  $u$  and  $v$  have continuous partial derivatives satisfying the Cauchy-Riemann conditions.*

$$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}; \quad \frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x} \quad (\text{C.6.1})$$

Furthermore

$$\frac{\partial w}{\partial z} = \frac{\partial u}{\partial x} + j\frac{\partial v}{\partial x} = \frac{\partial v}{\partial y} + j\frac{\partial v}{\partial x} = \frac{\partial u}{\partial x} - j\frac{\partial u}{\partial y} = \frac{\partial v}{\partial y} - j\frac{\partial u}{\partial y} \quad (\text{C.6.2})$$

**Proof**

Let  $z_0$  be a fixed point in  $D$  and let  $\Delta\omega = f(z_0 + \Delta z) - f(z_0)$ . Because  $f$  is analytic, we have

$$\Delta\omega = \gamma\Delta z + \epsilon\Delta z; \quad \gamma \triangleq f'(z_0) \quad (\text{C.6.3})$$

where  $\gamma = a + jb$  and  $\epsilon$  goes to zero as  $|z_0|$  goes to zero. Then

$$\Delta u + j\Delta v = (a + jb)(\Delta x + j\Delta y) + (\epsilon_1 + j\epsilon_2)(\Delta x + j\Delta y) \quad (\text{C.6.4})$$

So

$$\Delta u = a\Delta x - b\Delta y + \epsilon_1\Delta x - \epsilon_2\Delta y \quad (\text{C.6.5})$$

$$\Delta v = b\Delta x + a\Delta y + \epsilon_2\Delta x + \epsilon_1\Delta y \quad (\text{C.6.6})$$

Thus, in the limit, we can write

$$du = adx - bdy; \quad dv = bdx - ady \quad (\text{C.6.7})$$

or

$$\frac{\partial u}{\partial x} = a = -\frac{\partial v}{\partial y}; \quad \frac{\partial u}{\partial y} = -b = -\frac{\partial v}{\partial x} \quad (\text{C.6.8})$$

□□□

Actually, most functions that we will encounter will be analytic, provided the derivative exists. We illustrate this with some examples.

**Example C.1.** *Consider the function  $f(z) = z^2$ . Then*

$$f(z) = (x + jy)^2 = x^2 - y^2 + j(2xy) = u + jv \quad (\text{C.6.9})$$

The partial derivatives are

$$\frac{\partial u}{\partial x} = 2x; \quad \frac{\partial v}{\partial x} = 2y; \quad \frac{\partial u}{\partial y} = -2y; \quad \frac{\partial v}{\partial y} = 2x \quad (\text{C.6.10})$$

Hence, the function is clearly analytic.

**Example C.2.** Consider  $f(z) = |z|$ .

This function is not analytic, because  $d|z|$  is a real quantity and, hence,  $\frac{d|z|}{dz}$  will depend on the direction of  $z$ .

**Example C.3.** Consider a rational function of the form:

$$W(z) = K \frac{(z - \beta_1)(z - \beta_2) \cdots (z - \beta_m)}{(z - \alpha_1)(z - \alpha_2) \cdots (z - \alpha_n)} = \frac{N(z)}{D(z)} \quad (\text{C.6.11})$$

$$\frac{\partial W}{\partial z} = \frac{1}{D^2(z)} \left[ D(z) \frac{\partial N(z)}{\partial z} - N(z) \frac{\partial D(z)}{\partial z} \right] \quad (\text{C.6.12})$$

These derivatives clearly exist, save when  $D = 0$ , that is at the poles of  $W(z)$ .

**Example C.4.** Consider the same function  $W(z)$  defined in (C.6.11). Then

$$\frac{\partial \ln(W)}{\partial z} = \frac{1}{N(z)D(z)} \left[ D(z) \frac{\partial N(z)}{\partial z} - N(z) \frac{\partial D(z)}{\partial z} \right] = \frac{1}{N(z)} \frac{\partial N(z)}{\partial z} - \frac{1}{D(z)} \frac{\partial D(z)}{\partial z} \quad (\text{C.6.13})$$

Hence,  $\ln(W(z))$  is analytic, save at the poles and zeros of  $W(z)$ .

## C.7 Integrals Revisited

**Theorem C.7 (Cauchy Integral Theorem).** If  $f(z)$  is analytic in some simply connected domain  $D$ , then  $\int f(z)dz$  is independent of path in  $D$  and

$$\oint_C f(z)dz = 0 \quad (\text{C.7.1})$$

where  $C$  is a simple closed path in  $D$ .

### Proof

This follows from the Cauchy–Riemann conditions together with Theorem C.2.

□□□



We are also interested in the value of integrals in various limiting situations. The following examples cover relevant cases.

We note that if  $L_C$  is the length of a simple curve  $C$ , then

$$\left| \int_C f(z) dz \right| \leq \max_{z \in C} (|f(z)|) L_C \quad (\text{C.7.2})$$

**Example C.5.** Assume that  $C$  is a semicircle centered at the origin and having radius  $R$ . The path length is then  $L_C = \pi R$ . Hence,

- if  $f(z)$  varies as  $z^{-2}$ , then  $|f(z)|$  on  $C$  must vary as  $R^{-2}$  – hence, the integral on  $C$  vanishes for  $R \rightarrow \infty$ .
- if  $f(z)$  varies as  $z^{-1}$ , then  $|f(z)|$  on  $C$  must vary as  $R^{-1}$  – then, the integral on  $C$  becomes a constant as  $R \rightarrow \infty$ .

**Example C.6.** Consider the function  $f(z) = \ln(z)$  and an arc of a circle,  $C$ , described by  $z = \epsilon e^{j\gamma}$  for  $\gamma \in [-\gamma_1, \gamma_1]$ . Then

$$I_\epsilon \triangleq \lim_{\epsilon \rightarrow 0} \int_C f(z) dz = 0 \quad (\text{C.7.3})$$

This is proven as follows. On  $C$ , we have that  $f(z) = \ln(\epsilon)$ . Then

$$I_\epsilon = \lim_{\epsilon \rightarrow 0} [(\gamma_2 - \gamma_1)\epsilon \ln(\epsilon)] \quad (\text{C.7.4})$$

We then use the fact that  $\lim_{|x| \rightarrow 0} (x \ln x) = 0$ , and the result follows.

**Example C.7.** Consider the function

$$f(z) = \ln \left( 1 + \frac{a}{z^n} \right) \quad n \geq 1 \quad (\text{C.7.5})$$

and a semicircle,  $C$ , defined by  $z = Re^{j\gamma}$  for  $\gamma \in [-\frac{\pi}{2}, \frac{\pi}{2}]$ . Then, if  $C$  is followed clockwise,

$$I_R \triangleq \lim_{R \rightarrow \infty} \int_C f(z) dz = \begin{cases} 0 & \text{for } n > 1 \\ -j\pi a & \text{for } n = 1 \end{cases} \quad (\text{C.7.6})$$

This is proven as follows.

On  $C$ , we have that  $z = Re^{j\gamma}$ ; then

$$I_R = \lim_{R \rightarrow \infty} j \int_{\frac{\pi}{2}}^{-\frac{\pi}{2}} \ln \left( 1 + \frac{a}{R^n} e^{-jn\gamma} \right) Re^{j\gamma} d\gamma \quad (\text{C.7.7})$$

We also know that

$$\lim_{|x| \rightarrow 0} \ln(1+x) = x \quad (\text{C.7.8})$$

Then

$$I_R = \lim_{R \rightarrow \infty} \frac{a}{R^{n-1}} j \int_{\frac{\pi}{2}}^{-\frac{\pi}{2}} e^{-j(n-1)\gamma} d\gamma \quad (\text{C.7.9})$$

From this, by evaluation for  $n = 1$  and for  $n > 1$ , the result follows.

□□□

**Example C.8.** Consider the function

$$f(z) = \ln \left( 1 + e^{-z\tau} \frac{a}{z^n} \right) \quad n \geq 1; \quad \tau > 0 \quad (\text{C.7.10})$$

and a semicircle,  $C$ , defined by  $z = Re^{j\gamma}$  for  $\gamma \in [-\frac{\pi}{2}, \frac{\pi}{2}]$ . Then, for clockwise  $C$ ,

$$I_R \triangleq \lim_{R \rightarrow \infty} \int_C f(z) dz = 0 \quad (\text{C.7.11})$$

This is proven as follows.

On  $C$ , we have that  $z = Re^{j\gamma}$ ; then

$$I_R = \lim_{R \rightarrow \infty} j \int_{\frac{\pi}{2}}^{-\frac{\pi}{2}} \left[ \ln \left( 1 + \frac{a}{z(n+1)} \frac{z}{e^{z\tau}} \right) z \right]_{z=Re^{j\gamma}} d\gamma \quad (\text{C.7.12})$$

We recall that, if  $\tau$  is a positive real number and  $\Re\{z\} > 0$ , then

$$\lim_{|z| \rightarrow \infty} \frac{z}{e^{z\tau}} = 0 \quad (\text{C.7.13})$$

Moreover, for very large  $R$ , we have that

$$\ln \left( 1 + \frac{a}{z^{n+1}} \frac{z}{e^{z\tau}} \right) z \Big|_{z=Re^{j\gamma}} \approx \frac{1}{z^n} \frac{z}{e^{z\tau}} \Big|_{z=Re^{j\gamma}} \quad (\text{C.7.14})$$

Thus, in the limit, this quantity goes to zero for all positive  $n$ . The result then follows.

□□□

**Example C.9.** Consider the function

$$f(z) = \ln \left( \frac{z-a}{z+a} \right) \quad (\text{C.7.15})$$

and a semicircle,  $C$ , defined by  $z = Re^{j\gamma}$  for  $\gamma \in [-\frac{\pi}{2}, \frac{\pi}{2}]$ . Then, for clockwise  $C$ ,

$$I_R \triangleq \lim_{R \rightarrow \infty} \int_C f(z) dz = j2\pi a \quad (\text{C.7.16})$$

This result is obtained by noting that

$$\ln \left( \frac{z-a}{z+a} \right) = \ln \left( \frac{1 - \frac{a}{z}}{1 + \frac{a}{z}} \right) = \ln \left( 1 - \frac{a}{z} \right) - \ln \left( 1 + \frac{a}{z} \right) \quad (\text{C.7.17})$$

and then applying the result in example C.7.

□□□

**Example C.10.** Consider a function of the form

$$f(z) = \frac{a_{-1}}{z} + \frac{a_{-2}}{z^2} + \dots \quad (\text{C.7.18})$$

and  $C$ , an arc of circle  $z = Re^{j\theta}$  for  $\theta \in [\theta_1, \theta_2]$ . Thus,  $dz = jz d\theta$ , and

$$\int_C \frac{dz}{z} = \int_{\theta_1}^{\theta_2} j d\theta = -j(\theta_2 - \theta_1) \quad (\text{C.7.19})$$

Thus, as  $R \rightarrow \infty$ , we have that

$$\int_C f(z) dz = -ja_{-1}(\theta_2 - \theta_1) \quad (\text{C.7.20})$$

□□□

**Example C.11.** Consider, now,  $f(z) = z^n$ . If the path  $C$  is a full circle, centered at the origin and of radius  $R$ , then

$$\oint_C z^n dz = \int_{-\pi}^{\pi} (R^n e^{jn\theta}) j R e^{j\theta} d\theta \quad (\text{C.7.21})$$

$$= \begin{cases} 0 & \text{for } n \neq -1 \\ -2\pi j & \text{for } n = -1 \text{ (integration clockwise)} \end{cases} \quad (\text{C.7.22})$$

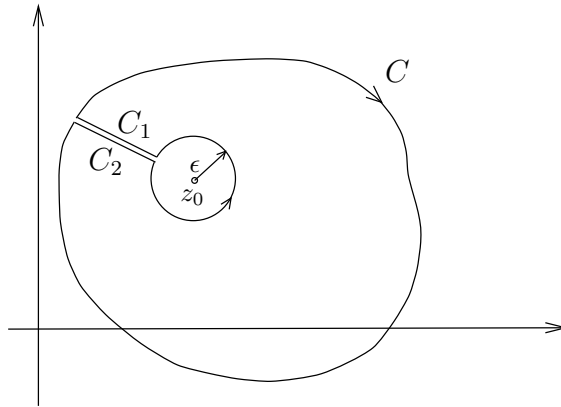
□□□

We can now develop Cauchy's Integral Formula.

Say that  $f(z)$  can be expanded as

$$f(z) = \frac{a_{-1}}{z - z_0} + a_0 + a_1(z - z_0) + a_2(z - z_0)^2 + \dots \quad (\text{C.7.23})$$

the  $a_{-1}$  is called the residue of  $f(z)$  at  $z_0$ .



**Figure C.3.** Path for integration of a function having a singularity

Consider the path shown in Figure C.3. Because  $f(z)$  is analytic in a region containing  $C$ , we have that the integral around the complete path shown in Figure C.3 is zero. The integrals along  $C_1$  and  $C_2$  cancel. The anticlockwise circular integral around  $z_0$  can be computed by following example C.11 to yield  $2\pi ja_{-1}$ . Hence, the integral around the outer curve  $C$  is minus the integral around the circle of radius  $\epsilon$ . Thus,

$$\oint_C f(z) dz = -2\pi ja_{-1} \quad (\text{C.7.24})$$

This leads to the following result.

**Theorem C.8 (Cauchy's Integral Formula).** *Let  $g(z)$  be analytic in a region. Let  $q$  be a point inside the region. Then  $\frac{g(z)}{z-q}$  has residue  $g(q)$  at  $z = q$ , and the integral around any closed contour  $C$  enclosing  $q$  in a clockwise direction is given by*

$$\oint_C \frac{g(z)}{z - q} dz = -2\pi jg(q) \quad (\text{C.7.25})$$

□□□

We note that the residue of  $g(z)$  at an interior point,  $z = q$ , of a region  $D$  can be obtained by integrating  $\frac{g(z)}{z-q}$  on the boundary of  $D$ . Hence, we can determine the value of an analytic function inside a region by its behaviour on the boundary.

## C.8 Poisson and Jensen Integral Formulas

We will next apply the Cauchy Integral formula to develop two related results.

The first result deals with functions that are analytic in the right-half plane (RHP). This is relevant to sensitivity functions in continuous-time systems, where Laplace transforms are used.

The second result deals with functions that are analytic outside the unit disk. This will be a preliminary step to analyzing sensitivity functions in discrete time, on the basis of Z-transforms.

### C.8.1 Poisson's Integral for the Half-Plane

**Theorem C.9.** Consider a contour  $C$  bounding a region  $D$ .  $C$  is a clockwise contour composed by the imaginary axis and a semicircle to the right, centered at the origin and having radius  $R \rightarrow \infty$ . This contour is shown in Figure C.4. Consider some  $z_0 = x_0 + jy_0$  with  $x_0 > 0$ .

Let  $f(z)$  be a real function of  $z$ , analytic inside  $D$  and of at least the order of  $z^{-1}$ ;  $f(z)$  satisfies

$$\lim_{|z| \rightarrow \infty} |z||f(z)| = \beta \quad 0 \leq \beta < \infty \quad z \in D \quad (\text{C.8.1})$$

then

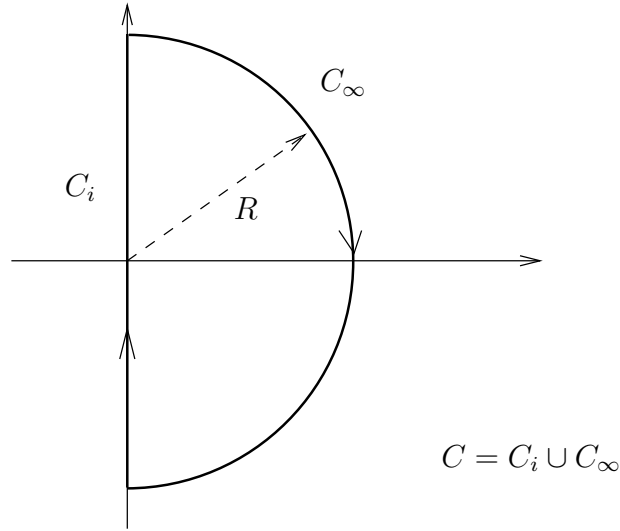
$$f(z_0) = -\frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{f(j\omega)}{j\omega - z_0} d\omega \quad (\text{C.8.2})$$

Moreover, if (C.8.1) is replaced by the weaker condition

$$\lim_{|z| \rightarrow \infty} \frac{|f(z)|}{|z|} = 0 \quad z \in D \quad (\text{C.8.3})$$

then

$$f(z_0) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(j\omega) \frac{x_0}{x_0^2 + (y_0 - \omega)^2} d\omega \quad (\text{C.8.4})$$



**Figure C.4.** RHP encircling contour

**Proof**

Applying Theorem C.8, we have

$$f(z_0) = -\frac{1}{2\pi j} \oint_C \frac{f(z)}{z - z_0} dz = -\frac{1}{2\pi j} \int_{C_i} \frac{f(z)}{z - z_0} dz - \frac{1}{2\pi j} \int_{C_\infty} \frac{f(z)}{z - z_0} dz \quad (\text{C.8.5})$$

Now, if  $f(z)$  satisfies (C.8.1), it behaves like  $z^{-1}$  for large  $|z|$ , i.e.,  $\frac{f(z)}{z - z_0}$  is like  $z^{-2}$ . The integral along  $C_\infty$  then vanishes and the result (C.8.2) follows.

To prove (C.8.4) when  $f(z)$  satisfies (C.8.3), we first consider  $z_1$ , the image of  $z_0$  through the imaginary axis, i.e.,  $z_1 = -x_0 + jy_0$ . Then  $\frac{f(z)}{z - z_1}$  is analytic inside  $D$ , and, on applying Theorem C.7, we have that

$$0 = -\frac{1}{2\pi j} \oint_C \frac{f(z)}{z - z_1} dz \quad (\text{C.8.6})$$

By combining equations (C.8.5) and (C.8.6), we obtain

$$f(z_0) = -\frac{1}{2j\pi} \oint_C \left( \frac{f(z)}{z - z_0} - \frac{f(z)}{z - z_1} \right) dz = -\frac{1}{2j\pi} \oint_C f(z) \frac{z_0 - z_1}{(z - z_0)(z - z_1)} dz \quad (\text{C.8.7})$$

Because  $C = C_i \cup C_\infty$ , the integral over  $C$  can be decomposed into the integral along the imaginary axis,  $C_i$ , and the integral along the semicircle of infinite radius,  $C_\infty$ . Because  $f(z)$  satisfies (C.8.3), this second integral vanishes, because the factor  $\frac{z_0 - z_1}{(z - z_0)(z - z_1)}$  is of order  $z^{-2}$  at  $\infty$ .

Then

$$f(z_0) = -\frac{1}{2\pi} \int_{-\infty}^{\infty} f(j\omega) \frac{z_0 - z_1}{(j\omega - z_0)(j\omega - z_1)} d\omega \quad (\text{C.8.8})$$

The result follows upon replacing  $z_0$  and  $z_1$  by their real; and imaginary-part decompositions. □□□

**Remark C.1.** One of the functions that satisfies (C.8.3) but does not satisfy (C.8.1) is  $f(z) = \ln g(z)$ , where  $g(z)$  is a rational function of relative degree  $n_r \neq 0$ . We notice that, in this case,

$$\lim_{|z| \rightarrow \infty} \left[ \frac{|\ln g(z)|}{|z|} \right] = \lim_{R \rightarrow \infty} \frac{|K| |n_r \ln R + j n_r \theta|}{R} = 0 \quad (\text{C.8.9})$$

where  $K$  is a finite constant and  $\theta$  is an angle in  $[-\frac{\pi}{2}, \frac{\pi}{2}]$ .

**Remark C.2.** Equation (C.8.4) equates two complex quantities. Thus, it also applies independently to their real and imaginary parts. In particular,

$$\Re\{f(z_0)\} = \frac{1}{\pi} \int_{-\infty}^{\infty} \Re\{f(j\omega)\} \frac{x_0}{x_0^2 + (y_0 - \omega)^2} d\omega \quad (\text{C.8.10})$$

This observation is relevant to many interesting cases. For instance, when  $f(z)$  is as in remark C.1,

$$\Re\{f(z)\} = \ln |g(z)| \quad (\text{C.8.11})$$

For this particular case, and assuming that  $g(z)$  is a real function of  $z$ , and that  $y_0 = 0$ , we have that (C.8.10) becomes

$$\ln |g(z_0)| = \frac{1}{\pi} \int_0^{\infty} \ln |g(j\omega)| \frac{2x_0}{x_0^2 + (y_0 - \omega)^2} d\omega \quad (\text{C.8.12})$$

where we have used the conjugate symmetry of  $g(z)$ .

### C.8.2 Poisson–Jensen Formula for the Half-Plane

**Lemma C.1.** Consider a function  $g(z)$  having the following properties

- (i)  $g(z)$  is analytic on the closed RHP;
- (ii)  $g(z)$  does not vanish on the imaginary axis;
- (iii)  $g(z)$  has zeros in the open RHP, located at  $a_1, a_2, \dots, a_n$ ;
- (iv)  $g(z)$  satisfies  $\lim_{|z| \rightarrow \infty} \frac{|\ln g(z)|}{|z|} = 0$ .

Consider also a point  $z_0 = x_0 + jy_0$  such that  $x_0 > 0$ ; then

$$\ln |g(z_0)| = \sum_{i=1}^n \ln \left| \frac{z_0 - a_i}{z_0 + a_i^*} \right| + \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{x_0}{x_0^2 + (\omega - y_0)^2} \ln |g(j\omega)| d\omega \quad (\text{C.8.13})$$

#### Proof

Let

$$\tilde{g}(z) \triangleq g(z) \prod_{i=1}^n \frac{z + a_i^*}{z - a_i} \quad (\text{C.8.14})$$

Then,  $\ln \tilde{g}(z)$  is analytic within the closed unit disk. If we now apply Theorem C.9 to  $\ln \tilde{g}(z)$ , we obtain

$$\ln \tilde{g}(z_0) = \ln g(z_0) + \sum_{i=1}^n \ln \left( \frac{z_0 + a_i^*}{z_0 - a_i} \right) = \frac{1}{\pi} \int_{-\infty}^{\infty} \frac{x_0}{x_0^2 + (\omega - y_0)^2} \ln \tilde{g}(j\omega) d\omega \quad (\text{C.8.15})$$

We also recall that, if  $x$  is any complex number, then  $\Re\{\ln x\} = \Re\{\ln |x| + j\angle x\} = \ln |x|$ . Thus, the result follows upon equating real parts in the equation above and noting that

$$\ln |\tilde{g}(j\omega)| = \ln |g(j\omega)| \quad (\text{C.8.16})$$

□□□

### C.8.3 Poisson's Integral for the Unit Disk

**Theorem C.10.** Let  $f(z)$  be analytic inside the unit disk. Then, if  $z_0 = re^{j\theta}$ , with  $0 \leq r < 1$ ,



$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} P_{1,r}(\theta - \omega) f(e^{j\omega}) d\omega \quad (\text{C.8.17})$$

where  $P_{1,r}(x)$  is the Poisson kernel defined by

$$P_{\rho,r}(x) \triangleq \frac{\rho^2 - r^2}{\rho^2 - 2r\rho \cos(x) + r^2} \quad 0 \leq r < \rho, \quad x \in \Re \quad (\text{C.8.18})$$

**Proof**

Consider the unit circle  $C$ . Then, using Theorem C.8, we have that

$$f(z_0) = \frac{1}{2\pi j} \oint_C \frac{f(z)}{z - z_0} dz \quad (\text{C.8.19})$$

Define

$$z_1 \triangleq \frac{1}{r} e^{j\theta} \quad (\text{C.8.20})$$

Because  $z_1$  is outside the region encircled by  $C$ , the application of Theorem C.8 yields

$$0 = \frac{1}{2\pi j} \oint_C \frac{f(z)}{z - z_1} dz \quad (\text{C.8.21})$$

Subtracting (C.8.21) from (C.8.19) and changing the variable of integration, we obtain

$$f(z_0) = \frac{1}{2\pi} \int_0^{2\pi} f(e^{j\omega}) e^{j\omega} \left[ \frac{1}{e^{j\omega} - r e^{j\theta}} - \frac{r}{r e^{j\omega} - e^{j\theta}} \right] d\omega \quad (\text{C.8.22})$$

from which the result follows. □□□

Consider now a function  $g(z)$  which is analytic outside the unit disk. We can then define a function  $f(z)$  such that

$$f(z) \triangleq g\left(\frac{1}{z}\right) \quad (\text{C.8.23})$$

Assume that one is interested in obtaining an expression for  $g(\zeta_0)$ , where  $\zeta_0 = re^{j\theta}$ ,  $r > 1$ . The problem is then to obtain an expression for  $f\left(\frac{1}{\zeta_0}\right)$ . Thus, if we define  $z_0 \triangleq \frac{1}{\zeta_0} = \frac{1}{r}e^{-j\theta}$ , we have, on applying Theorem C.10, that

$$g(\zeta_0) = \frac{1}{2\pi} \int_0^{2\pi} P_{1, \frac{1}{r}}(-\theta - \omega) g(e^{-j\omega}) d\omega \quad (\text{C.8.24})$$

where

$$P_{1, \frac{1}{r}}(-\theta - \omega) = \frac{r^2 - 1}{r^2 - 2r\cos(\theta + \omega) + 1} \quad (\text{C.8.25})$$

If, finally, we make the change in the integration variable  $\omega = -\nu$ , the following result is obtained.

$$g(re^{j\theta}) = \frac{1}{2\pi} \int_0^{2\pi} \frac{r^2 - 1}{r^2 - 2r\cos(\theta - \nu) + 1} g(e^{j\nu}) d\nu \quad (\text{C.8.26})$$

Thus, Poisson's integral for the unit disk can also be applied to functions of a complex variable which are analytic outside the unit circle.

#### C.8.4 Poisson–Jensen Formula for the Unit Disk

**Lemma C.2.** Consider a function  $g(z)$  having the following properties:

- (i)  $g(z)$  is analytic on the closed unit disk;
- (ii)  $g(z)$  does not vanish on the unit circle;
- (iii)  $g(z)$  has zeros in the open unit disk, located at  $\bar{\alpha}_1, \bar{\alpha}_2, \dots, \bar{\alpha}_n$ .

Consider also a point  $z_0 = re^{j\theta}$  such that  $r < 1$ ; then

$$\ln |g(z_0)| = \sum_{i=1}^n \ln \left| \frac{z_0 - \bar{\alpha}_i}{1 - \bar{\alpha}_i^* z_0} \right| + \frac{1}{2\pi} \int_0^{2\pi} P_{1,r}(\theta - \omega) \ln |g(e^{j\omega})| d\omega \quad (\text{C.8.27})$$

#### Proof

Let

$$\tilde{g}(z) \triangleq g(z) \prod_{i=1}^n \frac{1 - \bar{\alpha}_i^* z}{z - \bar{\alpha}_i} \quad (\text{C.8.28})$$

Then  $\ln \tilde{g}(z)$  is analytic on the closed unit disk. If we now apply Theorem C.10 to  $\ln \tilde{g}(z)$ , we obtain

$$\ln \tilde{g}(z_0) = \ln g(z_0) + \sum_{i=1}^n \ln \left( \frac{1 - \bar{\alpha}_i^* z_0}{z_0 - \bar{\alpha}_i} \right) = \frac{1}{2\pi} \int_0^{2\pi} P_{1,r}(\theta - \omega) \ln \tilde{g}(e^{j\omega}) d\omega \quad (\text{C.8.29})$$

We also recall that, if  $x$  is any complex number, then  $\ln x = \ln|x| + j\angle x$ . Thus the result follows upon equating real parts in the equation above and noting that

$$\ln |\tilde{g}(e^{j\omega})| = \ln |g(e^{j\omega})| \quad (\text{C.8.30})$$

□□□

**Theorem C.11 (Jensen's formula for the unit disk).** *Let  $f(z)$  and  $g(z)$  be analytic functions on the unit disk. Assume that the zeros of  $f(z)$  and  $g(z)$  on the unit disk are  $\bar{\alpha}_1, \bar{\alpha}_2, \dots, \bar{\alpha}_n$  and  $\bar{\beta}_1, \bar{\beta}_2, \dots, \bar{\beta}_m$  respectively, where none of these zeros lie on the unit circle.*

If

$$h(z) \triangleq z^\lambda \frac{f(z)}{g(z)} \quad \lambda \in \Re \quad (\text{C.8.31})$$

then

$$\frac{1}{2\pi} \int_0^{2\pi} \ln |h(e^{j\omega})| d\omega = \ln \left| \frac{f(0)}{g(0)} \right| + \ln \frac{|\bar{\beta}_1 \bar{\beta}_2 \dots \bar{\beta}_m|}{|\bar{\alpha}_1 \bar{\alpha}_2 \dots \bar{\alpha}_n|} \quad (\text{C.8.32})$$

**Proof**

We first note that  $\ln |h(z)| = \lambda \ln |z| + \ln |f(z)| - \ln |g(z)|$ . We then apply the Poisson–Jensen formula to  $f(z)$  and  $g(z)$  at  $z_0 = 0$  to obtain

$$P_{1,r}(x) = P_{1,0}(x) = 1; \quad \ln \left| \frac{z_0 - \bar{\alpha}_i}{1 - \bar{\alpha}_i^* z_0} \right| = \ln |\bar{\alpha}_i|; \quad \ln \left| \frac{z_0 - \bar{\beta}_i}{1 - \bar{\beta}_i^* z_0} \right| = \ln |\bar{\beta}_i| \quad (\text{C.8.33})$$

We thus have that

$$\ln |f(0)| = \sum_{i=1}^n \ln |\bar{\alpha}_i| - \frac{1}{2\pi} \int_0^{2\pi} \ln |f(e^{j\omega})| d\omega \quad (\text{C.8.34})$$

$$\ln |g(0)| = \sum_{i=1}^m \ln |\bar{\beta}_i| - \frac{1}{2\pi} \int_0^{2\pi} \ln |g(e^{j\omega})| d\omega \quad (\text{C.8.35})$$

The result follows upon subtracting equation (C.8.35) from (C.8.34), and noting that

$$\frac{\lambda}{2\pi} \int_0^{2\pi} \ln |e^{j\omega}| d\omega = 0 \quad (\text{C.8.36})$$

□□□

**Remark C.3.** Further insights can be obtained from equation (C.8.32) if we assume that, in (C.8.31),  $f(z)$  and  $g(z)$  are polynomials;

$$f(z) = K_f \prod_{i=1}^n (z - \alpha_i) \quad (\text{C.8.37})$$

$$g(z) = \prod_{i=1}^m (z - \beta_i) \quad (\text{C.8.38})$$

then

$$\left| \frac{f(0)}{g(0)} \right| = |K_f| \left| \frac{\prod_{i=1}^n \alpha_i}{\prod_{i=1}^m \beta_i} \right| \quad (\text{C.8.39})$$

Thus,  $\alpha_1, \alpha_2, \dots, \alpha_n$  and  $\beta_1, \beta_2, \dots, \beta_m$  are **all** the zeros and **all** the poles of  $h(z)$ , respectively, that have nonzero magnitude.

This allows equation (C.8.32) to be rewritten as

$$\frac{1}{2\pi} \int_0^{2\pi} \ln |h(e^{j\omega})| d\omega = \ln |K_f| + \ln \frac{|\alpha'_1 \alpha'_2 \dots \alpha'_{nu}|}{|\beta'_1 \beta'_2 \dots \beta'_{mu}|} \quad (\text{C.8.40})$$

where  $\alpha'_1, \alpha'_2, \dots, \alpha'_{nu}$  and  $\beta'_1, \beta'_2, \dots, \beta'_{mu}$  are the zeros and the poles of  $h(z)$ , respectively, that lie outside the unit circle.

□□□

### C.9 Application of the Poisson–Jensen Formula to Certain Rational Functions

Consider the biproper rational function  $\bar{h}(z)$  given by

$$\bar{h}(z) = z^{\bar{\lambda}} \frac{\bar{f}(z)}{\bar{g}(z)} \quad (\text{C.9.1})$$

$\bar{\lambda}$  is a integer number, and  $\bar{f}(z)$  and  $\bar{g}(z)$  are polynomials of degrees  $m_f$  and  $m_g$ , respectively. Then, due to the biproperness of  $\bar{h}(z)$ , we have that  $\bar{\lambda} + m_f = m_g$ .

Further assume that

- (i)  $\bar{g}(z)$  has no zeros outside the open unit disk,
- (ii)  $\bar{f}(z)$  does not vanish on the unit circle, and
- (iii)  $\bar{f}(z)$  vanishes outside the unit disk at  $\beta_1, \beta_2, \dots, \beta_m$ .

Define

$$h(z) = \frac{f(z)}{g(z)} \triangleq \bar{h}\left(\frac{1}{z}\right) \quad (\text{C.9.2})$$

where  $f(z)$  and  $g(z)$  are polynomials.

Then it follows that

- (i)  $g(z)$  has no zeros in the closed unit disk;
- (ii)  $f(z)$  does not vanish on the unit circle;
- (iii)  $f(z)$  vanishes in the open unit disk at  $\bar{\beta}_1, \bar{\beta}_2, \dots, \bar{\beta}_m$ , where  $\bar{\beta}_i = \beta_i^{-1}$  for  $i = 1, 2, \dots, \bar{\beta}_m$ ;
- (iv)  $h(z)$  is analytic in the closed unit disk;
- (v)  $h(z)$  does not vanish on the unit circle;
- (vi)  $h(z)$  has zeros in the open unit disk, located at  $\bar{\beta}_1, \bar{\beta}_2, \dots, \bar{\beta}_m$ .

We then have the following result

**Lemma C.3.** *Consider the function  $h(z)$  defined in (C.9.2) and a point  $z_0 = re^{j\theta}$  such that  $r < 1$ ; then*

$$\ln |h(z_0)| = \sum_{i=1}^{\bar{m}} \ln \left| \frac{z_0 - \bar{\beta}_i}{1 - \bar{\beta}_i^* z_0} \right| + \frac{1}{2\pi} \int_0^{2\pi} P_{1,r}(\theta - \omega) \ln |h(e^{j\omega})| d\omega \quad (\text{C.9.3})$$

where  $P_{1,r}$  is the Poisson kernel defined in (C.8.18).

**Proof**

This follows from a straightforward application of Lemma C.2.

□□□

## C.10 Bode's Theorems

We will next review some fundamental results due to Bode.

**Theorem C.12 (Bode integral in the half plane).** *Let  $l(z)$  be a proper real, rational function of relative degree  $n_r$ . Define*

$$g(z) \triangleq (1 + l(z))^{-1} \quad (\text{C.10.1})$$

and assume that  $g(z)$  has neither poles nor zeros in the closed RHP. Then

$$\int_0^\infty \ln |g(j\omega)| d\omega = \begin{cases} 0 & \text{for } n_r > 1 \\ -\kappa \frac{\pi}{2} & \text{for } n_r = 1 \end{cases} \quad \text{where } \kappa \triangleq \lim_{z \rightarrow \infty} z l(z) \quad (\text{C.10.2})$$

### Proof

Because  $\ln g(z)$  is analytic in the closed RHP,

$$\oint_C \ln g(z) dz = 0 \quad (\text{C.10.3})$$

where  $C = C_i \cup C_\infty$  is the contour defined in Figure C.4.

Then

$$\oint_C \ln g(z) dz = j \int_{-\infty}^\infty \ln g(j\omega) d\omega - \int_{C_\infty} \ln(1 + l(z)) dz \quad (\text{C.10.4})$$

For the first integral on the right-hand side of equation (C.10.4), we use the conjugate symmetry of  $g(z)$  to obtain

$$\int_{-\infty}^\infty \ln g(j\omega) d\omega = 2 \int_0^\infty \ln |g(j\omega)| d\omega \quad (\text{C.10.5})$$

For the second integral, we notice that, on  $C_\infty$ ,  $l(z)$  can be approximated by

$$\frac{a}{z^{n_r}} \quad (\text{C.10.6})$$

The result follows upon using example C.7 and noticing that  $a = \kappa$  for  $n_r = 1$ .  
□□□

**Remark C.4.** *If  $g(z) = (1 + e^{-z\tau} l(z))^{-1}$  for  $\tau > 0$ , then result (C.10.9) becomes*

$$\int_0^\infty \ln |g(j\omega)| d\omega = 0 \quad \forall n_r > 0 \quad (\text{C.10.7})$$

The proof of (C.10.7) follows along the same lines as those of Theorem C.12 and by using the result in example C.8.

**Theorem C.13 (Modified Bode integral).** Let  $l(z)$  be a proper real, rational function of relative degree  $n_r$ . Define

$$g(z) \triangleq (1 + l(z))^{-1} \quad (\text{C.10.8})$$

Assume that  $g(z)$  is analytic in the closed RHP and that it has  $q$  zeros in the open RHP, located at  $\zeta_1, \zeta_2, \dots, \zeta_q$  with  $\Re(\zeta_i) > 0$ . Then

$$\int_0^\infty \ln |g(j\omega)| d\omega = \begin{cases} \pi \sum_{i=1}^q \zeta_i & \text{for } n_r > 1 \\ -\kappa \frac{\pi}{2} + \pi \sum_{i=1}^q \zeta_i & \text{for } n_r = 1 \end{cases} \quad \text{where } \kappa \triangleq \lim_{z \rightarrow \infty} z l(z) \quad (\text{C.10.9})$$

### Proof

We first notice that  $\ln g(z)$  is no longer analytic on the RHP. We then define

$$\tilde{g}(z) \triangleq g(z) \prod_{i=1}^q \frac{z + \zeta_i}{z - \zeta_i} \quad (\text{C.10.10})$$

Thus,  $\ln \tilde{g}(z)$  is analytic in the closed RHP. We can then apply Cauchy's integral in the contour  $C$  described in Figure C.4 to obtain

$$\oint_C \ln \tilde{g}(z) dz = 0 = \oint_C \ln g(z) dz + \sum_{i=1}^q \oint_C \ln \frac{z + \zeta_i}{z - \zeta_i} dz \quad (\text{C.10.11})$$

The first integral on the right-hand side can be expressed as

$$\oint_C \ln g(z) dz = 2j \int_0^\infty \ln |g(j\omega)| d\omega + \int_{C_\infty} \ln g(z) dz \quad (\text{C.10.12})$$

where, by using example C.7.

$$\int_{C_\infty} \ln g(z) dz = \begin{cases} 0 & \text{for } n_r > 1 \\ j\kappa\pi & \text{for } n_r = 1 \end{cases} \quad \text{where } \kappa \triangleq \lim_{z \rightarrow \infty} z l(z) \quad (\text{C.10.13})$$



The second integral on the right-hand side of equation (C.10.11) can be computed as follows:

$$\oint_C \ln \frac{z + \zeta_i}{z - \zeta_i} dz = j \int_{-\infty}^{\infty} \ln \frac{j\omega + \zeta_i}{j\omega - \zeta_i} d\omega + \int_{C_\infty} \ln \frac{z + \zeta_i}{z - \zeta_i} dz \quad (\text{C.10.14})$$

We note that the first integral on the right-hand side is zero, and by using example C.9, the second integral is equal to  $-2j\pi\zeta_i$ . Thus, the result follows.  $\square\square\square$

**Remark C.5.** Note that  $g(z)$  is a real function of  $z$ , so

$$\sum_{i=1}^q \zeta_i = \sum_{i=1}^q \Re\{\zeta_i\} \quad (\text{C.10.15})$$

$\square\square\square$

**Remark C.6.** If  $g(z) = (1 + e^{-z\tau}l(z))^{-1}$  for  $\tau > 0$ , then the result (C.10.9) becomes

$$\int_0^{\infty} \ln |g(j\omega)| d\omega = \pi \sum_{i=1}^q \Re\{\zeta_i\} \quad \forall n_r > 0 \quad (\text{C.10.16})$$

The proof of (C.10.16) follows along the same lines as those of Theorem C.13 and by using the result in example C.8.

**Remark C.7.** The Poisson, Jensen, and Bode formulae assume that a key function is analytic, not only inside a domain  $D$ , but also on its border  $C$ . Sometimes, there may exist singularities on  $C$ . These can be dealt with by using an infinitesimal circular indentation in  $C$ , constructed so as to leave the singularity outside  $D$ . For the functions of interest to us, the integral along the indentation vanishes. This is illustrated in example C.6 for a logarithmic function, when  $D$  is the right-half plane and there is a singularity at the origin.

$\square\square\square$